

# Dirac Operators for Kinetic and Potential Energy

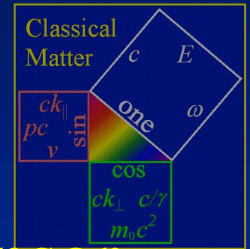
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# Spin Density



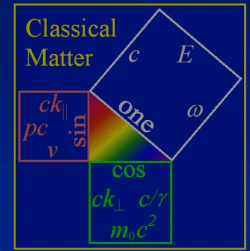
Belinfante-Rosenfeld symmetric stress-energy tensor implies dynamic (wave) & conjugate momenta:

$$\begin{aligned} \mathbf{P}_{Total} &= -\text{Re}[\psi^\dagger i\nabla\psi] + \frac{1}{2} \nabla \times \left[ \psi^\dagger \frac{\boldsymbol{\sigma}}{2} \psi \right] \\ &= \mathbf{P}_{wave} + \frac{1}{2} \nabla \times \mathbf{s} \end{aligned}$$

$$\begin{aligned} \mathbf{J} &= -\text{Re}(\mathbf{r} \times [\psi^\dagger i\nabla\psi]) + \left[ \psi^\dagger \frac{\boldsymbol{\sigma}}{2} \psi \right] \\ &= \mathbf{L}_{wave} + \mathbf{s} \end{aligned}$$

Elastic waves in a solid also have dynamic & conjugate (or “intrinsic”) momenta.

# Classical Spin Density



Incompressible Helmholtz decomposition:

$$\mathbf{p} \equiv \rho \mathbf{u} = \mathbf{p}_0 + \nabla \Phi + \frac{1}{2} \nabla \times \mathbf{s}$$

Angular momentum:

$$\mathbf{S} = \int \mathbf{r} \times \mathbf{p} d^3r = \int \mathbf{s} d^3r + \text{b.t.}$$

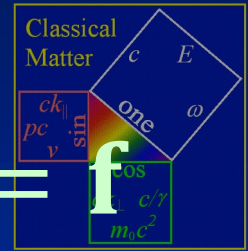
Kinetic energy:  $\left( \mathbf{w} = \frac{1}{2} \nabla \times \mathbf{u} \right)$

$$K = \int \frac{1}{2} \rho u^2 d^3r = \int \frac{1}{2} \mathbf{w} \cdot \mathbf{s} d^3r + \text{b.t.}$$

Momentum conjugate to  $\mathbf{w}$ : ( $\mathcal{L} \sim K$ )

$$\frac{\partial \mathcal{L}}{\partial \mathbf{w}} = \frac{\mathbf{s}}{2} + \frac{\mathbf{s}}{2} = \mathbf{s}$$

# Equation of Evolution



Momentum density:  $\partial_t \mathbf{p} + \mathbf{u} \cdot \nabla \mathbf{p} = \mathbf{f}$

Changes due to convection & force.

Spin density:  $\partial_t \mathbf{s} + \mathbf{u} \cdot \nabla \mathbf{s} - \mathbf{w} \times \mathbf{s} = \boldsymbol{\tau}$

Changes  $\sim$  convection, rotation, & torque.

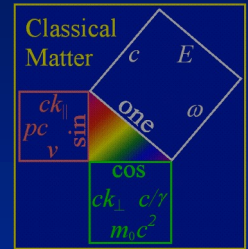
$\mathbf{s} \equiv \partial_t \mathbf{Q}$ , Displacement:  $\boldsymbol{\xi} = \frac{1}{2\rho} \nabla \times \mathbf{Q}$

Elastic solid: Torque =  $\boldsymbol{\tau} = -\frac{\mu}{\rho} \nabla \times \nabla \times \mathbf{Q}$

$\partial_t^2 \mathbf{Q} + \mathbf{u} \cdot \nabla \partial_t \mathbf{Q} - \mathbf{w} \times \partial_t \mathbf{Q} = c^2 \nabla^2 \mathbf{Q}$

Nonlinearity  $\Rightarrow$  soliton solutions

# Dirac Equation



Linear vector wave equations:

$$0 = \partial_t [\psi^\dagger \boldsymbol{\sigma} \psi] + c \nabla [\psi^\dagger \gamma^5 \psi] \\ - ic [\nabla \psi^\dagger \times \gamma^5 \boldsymbol{\sigma} \psi + \psi^\dagger \gamma^5 \boldsymbol{\sigma} \times \nabla \psi]$$

$$0 = \partial_t^2 \mathbf{Q} - c^2 \nabla (\nabla \cdot \mathbf{Q}) + c^2 \nabla \times \nabla \times \mathbf{Q} \\ = \partial_t^2 \mathbf{Q} - c^2 \nabla^2 \mathbf{Q}$$

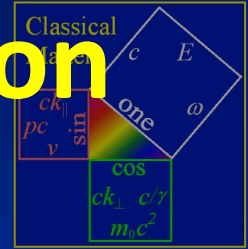

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$$\mathbf{s} = \partial_t \mathbf{Q} = (1/2) \psi^\dagger \boldsymbol{\sigma} \psi$$

$$c \nabla \cdot \mathbf{Q} = -(1/2) \psi^\dagger \gamma^5 \psi$$

$$c \nabla \times \nabla \times \mathbf{Q} = -\frac{i}{2} [\nabla \psi^\dagger \times \gamma^5 \boldsymbol{\sigma} \psi + \psi^\dagger \gamma^5 \boldsymbol{\sigma} \times \nabla \psi]$$

# Nonlinear Equation of Evolution



Nonlinear Dirac eqn.  $\mathbf{s} \equiv \partial_t \mathbf{Q} = \left[ \psi^\dagger \frac{\boldsymbol{\sigma}}{2} \psi \right]:$

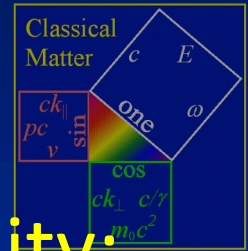
$\partial_t \psi =$

$$-c \gamma^5 \boldsymbol{\sigma} \cdot \nabla \psi - \mathbf{u} \cdot \nabla \psi - \frac{i}{2} \check{\omega}_0 \gamma^0 \psi - \frac{i}{2} \mathbf{w} \cdot \boldsymbol{\sigma} \psi$$

Propagation, convection, rotation of wave velocity, rotation of the solid medium.

Mass represents rotation of wave velocity.\*

\*c.f. Hestenes, *Found. Phys.* **20**(10):1213-32,1990



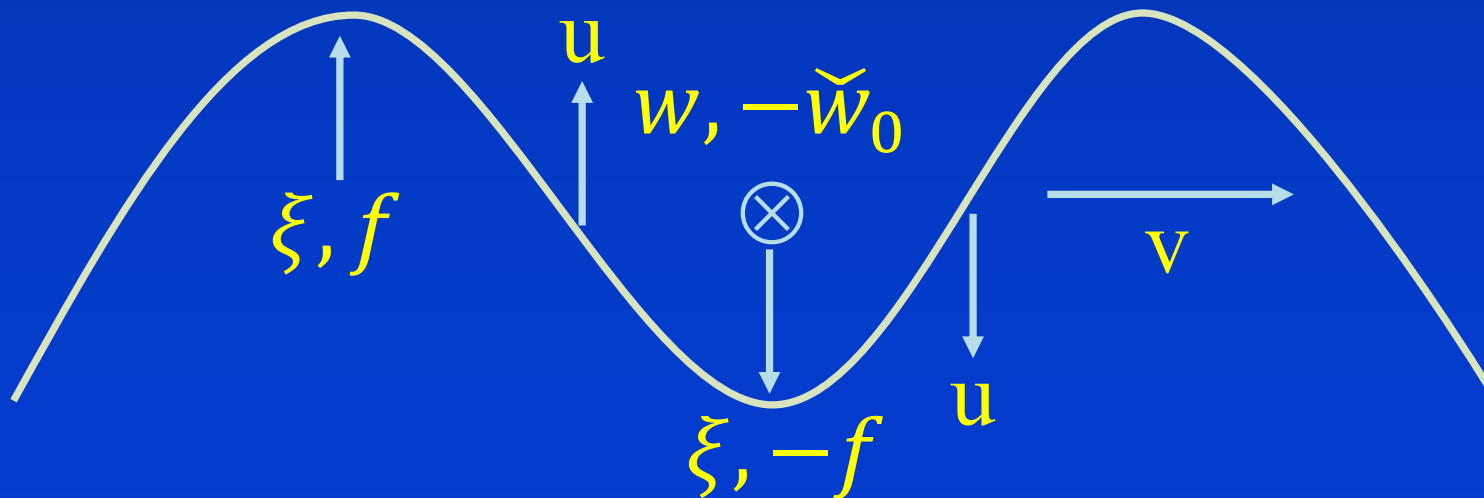
# Potential Energy

$\psi^\dagger \gamma^0 \psi \sim$  displacement.  $\check{w}_0 \sim$  force density:

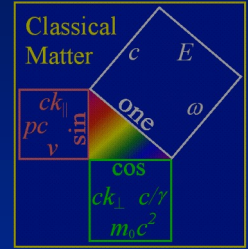
$$\check{w}_0 = \frac{1}{2\rho c} f = \frac{1}{4\rho} \nabla^2 \left( \psi^\dagger \frac{\gamma^0}{2} \psi \right)$$

PE density is:  $U = \frac{\check{w}_0}{4} \psi^\dagger \gamma^0 \psi \sim -\frac{1}{2} \mathbf{f} \cdot \boldsymbol{\xi}$

Rotational PE:  $U_R = \mathcal{E} - U$



# Hamiltonian Density



$$\begin{aligned}
 \mathcal{H} &= \text{Re}(-c\psi^\dagger \gamma^5 \boldsymbol{\sigma} \cdot i\nabla\psi - \mathbf{u} \cdot \psi^\dagger i\nabla\psi) \\
 &\quad + \frac{\check{W}_0}{4} \psi^\dagger \gamma^0 \psi + \frac{\mathbf{w}}{4} \cdot \psi^\dagger \boldsymbol{\sigma} \psi \\
 &= \left( \varepsilon - 0 + \frac{1}{2} \mathbf{f} \cdot \boldsymbol{\xi} \right) + \frac{1}{2} \mathbf{w} \cdot \mathbf{s} \\
 &= U_R + K_R
 \end{aligned}$$

Rotational potential energy:  $U_R = \varepsilon + \frac{1}{2} \mathbf{f} \cdot \boldsymbol{\xi}$

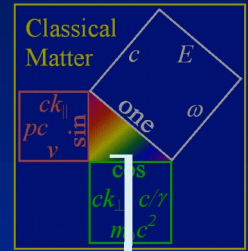
$$= \text{Re}(-c\psi^\dagger \gamma^5 \boldsymbol{\sigma} \cdot i\nabla\psi) + \frac{\check{W}_0}{4} \psi^\dagger \gamma^0 \psi$$

Rotational kinetic energy:  $K_R = \frac{1}{2} \mathbf{w} \cdot \mathbf{s} = \varepsilon - \frac{1}{2} \rho u^2$

$$= \frac{1}{4} \mathbf{w} \cdot \psi^\dagger \boldsymbol{\sigma} \psi$$



# Example: Plane Wave



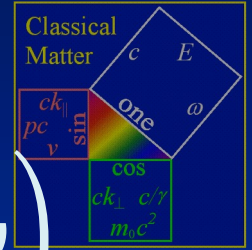
$$\psi = \sqrt{\frac{\omega Q_0}{2}} \begin{bmatrix} 1 \\ \cos(\omega t - kz) - i \frac{\omega Q_0 k^2}{4\rho c} \sin(\omega t - kz) \\ \cos(\omega t - kz) - i \frac{\omega Q_0 k^2}{4\rho c} \sin(\omega t - kz) \\ 1 \end{bmatrix}$$

$$s_x = \frac{1}{2} \psi^\dagger \sigma_x \psi = \omega Q_0 \cos(\omega t - kz)$$

$$s_y = \frac{1}{2} \psi^\dagger \sigma_y \psi = 0$$

$$s_z = \frac{1}{2} \psi^\dagger \sigma_z \psi = 0$$

# Example: Plane Wave



$$U_R = \left( -\text{Re}(\psi^\dagger \gamma^5 \boldsymbol{\sigma} \cdot i\nabla \psi) + \frac{1}{4} \check{w}_0 \psi^\dagger \gamma^0 \psi \right)$$

$$= \frac{k^2 \omega^2 Q_0^2}{8\rho} \sin^2(\omega t - kz)$$

$$K_R = \frac{1}{4} \mathbf{w} \cdot \boldsymbol{\sigma} \psi = \frac{k^2 \omega^2 Q_0^2}{8\rho} \cos^2(\omega t - kz)$$

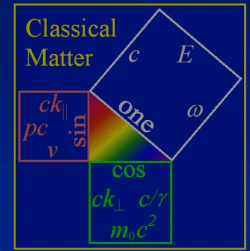
$$i\partial_t \psi = \left( -\gamma^5 \boldsymbol{\sigma} \cdot i\nabla \psi + \frac{1}{2} \check{w}_0 \gamma^0 \psi \right) + \frac{1}{2} \mathbf{w} \cdot \boldsymbol{\sigma} \psi$$

$$\mathcal{E} = \mathbf{v} \cdot \mathbf{P} - 2U + 2K_R$$

RHS  $\neq$  Hamiltonian (because nonlinear).

QM sets  $\mathbf{v} \cdot \mathbf{P} = 0$  and  $m_e c^2 = 2U$ .

# Operators as Generators

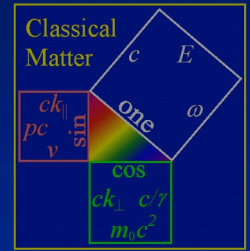


$$\mathbf{P}\psi = -i\nabla\psi + \frac{1}{2}\nabla \times \frac{\boldsymbol{\sigma}}{2}\psi$$

$$\mathbf{J}\psi = -\mathbf{r} \times i\nabla\psi + \frac{\boldsymbol{\sigma}}{2}\psi$$

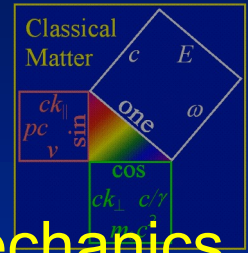
- $\mathbf{J}$  is generator of rotations (orbital for coordinate, spin for direction)
- $\mathbf{P}$  is generator of translations.
  - Wave momentum shifts coordinate ( $\mathbf{r}$ ).
  - Intrinsic momentum shifts displacement from equilibrium ( $\boldsymbol{\xi}$ ).

# Summary



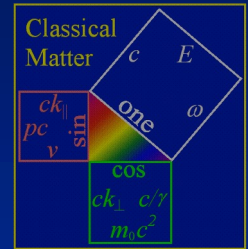
- Dirac equation describes spin density.
- Spin is ordinary angular momentum due to rotational motion of an inertial medium.
- Elastic solid model of vacuum yields nonlinear Dirac eqn. with Hamiltonian = Total energy
- Kinetic energy  $\sim$  angular velocity  $\cdot$  spin density.
- Potential energy  $\sim$  force density  $\cdot$  displacement  
(force density  $\sim$  wave velocity rotation rate)
- Solid model also yields exclusion principle, interaction potentials, fermions, bosons, etc.

# Publications



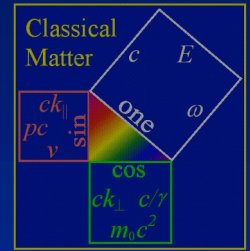
- Torsion Waves in Three Dimensions: Quantum Mechanics With a Twist," **Found. Phys. Lett.** 15(1):71-83, Feb. 2002.
- "A Classical Dirac Bispinor Equation," in *Ether Space-time & Cosmology*, vol. 3, eds. M. C. Duffy & J. Levy, (Aperion, Montreal, 2009).
- "Exact Description of Rotational Waves in an Elastic Solid," **Adv. App. Clifford Algebras** 21:273-281, 2010.
- *The Wave Basis of Special Relativity* (Verum Versa, 2014)
- "Spin Angular Momentum and the Dirac Equation," **Electr. J. Theor. Phys.** 12(33):43-60, 2015. url = <https://www.arcane-editrice.it/index.php/pubblicazione.html?item=9788854889095>
- More at: [www.ClassicalMatter.org](http://www.ClassicalMatter.org)

# Related Publications



- de Broglie LV 1924 *Recherches sur la Theorie des Quanta*, PhD Thesis, (Paris: University of Sorbonne).
- Einstein A 1956 *The Meaning of Relativity* (Princeton: Princeton Univ. Press) Fifth Edition, pp 84-89.
- Evans JC et. al. 2001 Matter waves in a gravitational field: An index of refraction for massive particles in general relativity, *Am. J. Phys.* **69**, 1103--10.
- Gu YQ 1998 Some Properties of the Spinor Soliton, *Advances in Applied Clifford Algebras* **8**(1) 17-29.
- Hestenes D 1990 The Zitterbewegung Interpretation of Quantum Mechanics, *Found. Phys.* **20**(10) 1213-32.
- Jehle H 1971, Relationship of Flux Quantization to Charge Quantization and the Electromagnetic Coupling Constant, *Phys. Rev.* **3**(2):306-345.
- Karlsten BU 1998 Sketch of a Matter Model in an Elastic Universe (<http://home.online.no/~ukarlsten>).
- Kleinert H 1989 *Gauge Fields in Condensed Matter* vol II (Singapore: World Scientific) pp 1259.
- Laughlin R. 2005, *A Different Universe* (New York, Basic Books).
- Lee TD and Yang CN 1956 Question of Parity Conservation in Weak Interactions, *Phys. Rev.* **104**, 254.
- Morse PM and Feshbach H 1953a *Methods of Theoretical Physics* vol I (New York: McGraw-Hill Book Co.) pp 304-6.
- Ohanian HC 1986 What is Spin, *Am. J. Phys.* **54**(6):500-5.
- Ranada AF 1983 Classical Nonlinear Dirac Field Models of Extended Particles *Quantum Theory, Groups, Fields, and Particles* ed A O Barut (Amsterdam: Reidel) pp 271-88.
- Rowlands P 1998 The physical consequences of a new version of the Dirac equation *Causality and Locality in Modern Physics and Astronomy: Open Questions and Possible Solutions* (Fundamental Theories of Physics, vol 97) eds G. Hunter, S. Jeffers, and J-P. Vigiier (Dordrecht: Kluwer Academic Publishers) pp 397-402.
- Rowlands P 2005 Removing redundancy in relativistic quantum mechanics *Preprint* arXiv:physics/0507188.
- Schmeltzer I 2012 The standard model fermions as excitations of an ether, in *Horizons in World Physics, vol. 278*, edited by A. Reimer, (Nova Science Publishers).
- Steinberg DJ, Cochran SG, & Guinana MW 1980 A constitutive model for metals applicable at high-strain rate, *J. Appl. Phys.* **51**:1498-1504.
- Takabayashi Y 1957 Relativistic hydrodynamics of the Dirac matter, *Suppl. Prog. Theor. Phys.* **4**(1) 1-80.
- Whittaker E 1951 *A History of the Theories of Aether and Electricity*, (Edinburgh: Thomas Nelson and Sons Ltd.).
- Wilson HA 1921 An Electromagnetic Theory of Gravitation, *Phys. Rev.* **17**: 54-59.
- Yamamoto H 1977 Spinor soliton as an elementary particle, *Prog. Theor. Phys.* **58**(3), 1014--23 .

# Conservation Law



Nonlinear Dirac equation:

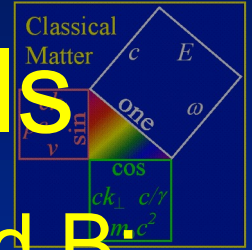
$$\partial_t \psi = -c \gamma^5 \boldsymbol{\sigma} \cdot \nabla \psi - \mathbf{u} \cdot \nabla \psi - \frac{i}{2} (\tilde{w}_0 \gamma^0 \psi + \mathbf{w} \cdot \boldsymbol{\sigma} \psi)$$

Multiply  $\psi^\dagger$  and add adjoint:

$$\partial_t (\psi^\dagger \psi) = -\nabla \cdot (\psi^\dagger c \gamma^5 \boldsymbol{\sigma} \psi) - \mathbf{u} \cdot \nabla (\psi^\dagger \psi)$$

Magnitude  $\psi^\dagger \psi$  is a conserved quantity.

# Exclusion Principle & Potentials



- Wave superposition of “particles” A and B:

$$[\psi_A + \psi_B]^\dagger [\psi_A + \psi_B] = \psi_A^\dagger \psi_A + \psi_B^\dagger \psi_B + \psi_A^\dagger \psi_B + \psi_B^\dagger \psi_A$$

- Interference terms cancel for “independent” particles  $\Rightarrow$  exclusion principle:

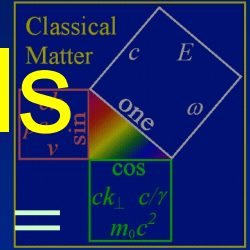
$$\psi_A^\dagger \psi_B + \psi_B^\dagger \psi_A = 0$$

- Potentials derive from phase shifts introduced to maintain zero interference.

$$\psi_A^\dagger e^{-i\varphi_A} e^{i\varphi_B} \psi_B + c.c. = 0$$



# Exclusion Principle & Potentials



Change of wave momentum (force):  $\frac{d}{dt} P_i =$

$$\psi_A^\dagger \left( (\partial_t \partial_i - \partial_i \partial_t) \varphi_A - c \gamma^5 \sigma_j (\partial_i \partial_j - \partial_j \partial_i) \varphi_A \right) \psi_A$$

Multivalued phase  $\Rightarrow$  Non-commuting!

Electromagnetic variables:

$$e\mathbf{A} = -\hbar \nabla \varphi_A; \quad e\Phi = \hbar \partial_t \varphi_A$$

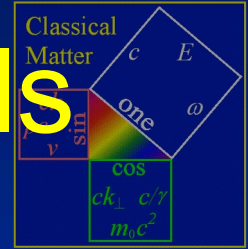
$$\rho_e = e \psi_A^\dagger \psi_A; \quad \mathbf{J} = e \psi_A^\dagger c \gamma^5 \boldsymbol{\sigma} \psi_A$$

$$\mathbf{E} = -\frac{\hbar}{e} \nabla (\partial_t \varphi_A) + \frac{\hbar}{e} \partial_t (\nabla \varphi_A) = -\nabla \Phi - \partial_t \mathbf{A}$$

$$B_k = \epsilon_{kij} \partial_i A_j = -\frac{\hbar}{e} \epsilon_{kij} \partial_i \partial_j \varphi_A$$

$$\frac{d}{dt} \mathbf{P} = \rho_e \mathbf{E} + \mathbf{J} \times \mathbf{B}$$

# Exclusion Principle & Potentials



Stokes' Law:  $\oint \mathbf{A} \cdot d\boldsymbol{\ell} = \iint (\nabla \times \mathbf{A}) \cdot d\mathbf{S}$

Macroscopic system:

Pure phase shift  $\Rightarrow$  magnetic flux quantized:

$$\varphi_A = (m_\phi \phi - \omega t)$$

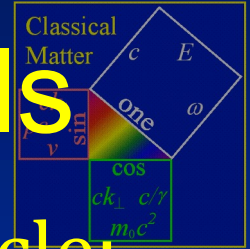
$$\mathbf{A} = -\frac{\hbar}{e} \nabla \varphi_A = -\frac{\hbar m_\phi}{e r \sin \theta}$$

$$\mathbf{B} = \nabla \times \mathbf{A} = \frac{\hbar m_\phi}{e} 2\pi \delta^{(2)}(x, y)$$

$$\oint \mathbf{A} \cdot d\boldsymbol{\ell} = -\frac{\hbar}{e} 2\pi m_\phi \xrightarrow{m_\phi=1/2} \frac{h}{2e}$$

c.f. Hagen Kleinert, *Multivalued Fields in Condensed Matter, Electromagnetism, and Gravitation* (Chapter 4).

# Exclusion Principle & Potentials



Radially weighted phase shift from particle:

$$\varphi_A = (m_\phi \phi - \omega t) \left( \frac{e^2}{4\pi\epsilon_0 \hbar c} \right) \frac{c}{\omega r}$$

$$e\mathbf{A} = \left( \frac{e^2}{4\pi\epsilon_0 \hbar c} \right) \frac{c}{\omega r} \left\{ \left( \frac{m_\phi}{r \sin \theta} \right) \hat{\phi} - \frac{m_\phi \phi - \omega t}{r^2} \hat{r} \right\}$$

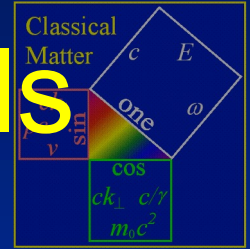
Neglect radial term (no flux). Note  $\hbar\omega = m_e c^2$

$$e\mathbf{B} = \left( \frac{e^2}{4\pi\epsilon_0 \hbar c} \right) \frac{c}{\omega} \left( \frac{m_\phi}{r^3 \sin \theta} \right) \hat{\theta}$$

Same flux as  $e^-$  dipole moment:  $\mu_0 M / 2r$ :

$$\oint \mathbf{A} \cdot d\boldsymbol{\ell} = -\frac{m_\phi e}{2\epsilon_0 \omega r} \xrightarrow{m_\phi=1/2} \frac{\mu_0 \hbar e}{4m_e r}$$

# Exclusion Principle & Potentials



Radially weighted phase shift:

$$\varphi_A = (m_\phi \phi - \omega t) \left( \frac{e^2}{4\pi\epsilon_0 \hbar c} \right) \frac{c}{\omega r}$$

Phase velocity:  $v_\phi = \frac{\omega}{m_\phi} r \sin \theta \rightarrow 2\omega r \sin \theta$

Electric field (Note  $\partial_t \varphi_A = -\mathbf{v} \cdot \nabla \varphi_A$ ):

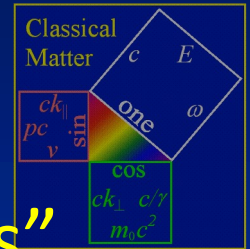
$$eE_i = \hbar(\partial_i v_j) \partial_j \varphi_A + \hbar v_j (\partial_i \partial_j - \partial_j \partial_i) \varphi_A$$

Only first term contributes ( $v_\phi = 0$  at  $r = 0$ )

$$\mathbf{E} = \left( \frac{e}{4\pi\epsilon_0 r^2} \right) \hat{\mathbf{r}}$$

*c.f.* Herbert Jehle, Phys. Rev. 3(2):306-345, 1971

# Matter & Antimatter



Assume vector spherical harmonic “particles”.

Bispinor angular quantum number is half of vector angular quantum number.

Vector:  $\ell = 2N + 1 \Rightarrow$  Bispinor:  $\ell = \frac{2N+1}{2}$

Odd parity  $\Rightarrow$  Distinct mirror image  $\Rightarrow$  Fermion

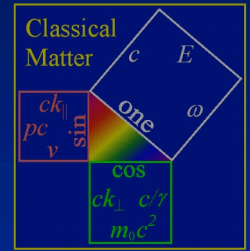
Vector:  $\ell = 2N \Rightarrow$  Bispinor:  $\ell = N$

Even parity  $\Rightarrow$  Same mirror image  $\Rightarrow$  Boson

(bosons = antiparticles except  $W+/-$ )

Assumption valid except for  $W+/-$ .

# Quantum Electrodynamics



2-Particle Lagrangian:

$$\mathcal{L} = \bar{\psi}_A [\gamma^\mu (i \partial_\mu - A_\mu) - m_A] \psi_A + (A \rightarrow B)$$

Apply Dirac equation to B:  $\mathcal{L} =$

$$\bar{\psi}_A [\gamma^\mu (i \partial_\mu - A_\mu) - m_A] \psi_A + J^\mu A_\mu - J^\mu A_\mu$$

Green's 1<sup>st</sup> identity, vector identities, etc.:

$$\int J^\mu A_\mu = \int (E^2 - B^2) = -\frac{1}{2} \int F^{\mu\nu} F_{\mu\nu}$$

Factor of 1/2 due to counting variations twice:

$$\begin{aligned} \mathcal{L} = & \bar{\psi}_A [\gamma^\mu (i \partial_\mu - A_\mu) - m_A] \psi_A \\ & + \frac{1}{2} (E^2 - B^2) - J^\mu A_\mu \end{aligned}$$